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Technical Note: A novel parameterization of the transmissivity due to ozone absorption in the *k*-distribution method and correlated-*k* approximation of Kato et al. (1999) over the UV band

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Abstract. The *k*-distribution method and the correlated-*k* approximation of Kato et al. (1999) is a computationally efficient approach originally designed for calculations of the broadband solar radiation at ground level by dividing the solar spectrum in 32 specific spectral bands from 240 to 4606 nm. Compared to a spectrally resolved computation, its performance in the UV band appears to be inaccurate, especially in the spectral intervals #3 [283, 307] nm and #4 [307, 328] nm because of inaccuracy in modeling the transmissivity due to ozone absorption. Numerical simulations presented in this paper indicate that a single effective ozone cross section is insufficient to accurately represent the transmissivity over each spectral interval. A novel parameterization of the transmissivity using more quadrature points yields maximum errors of respectively 0.0006 and 0.0143 for intervals #3 and #4. How to practically implement this new parameterization in a radiative transfer model is discussed for the case of libRadtran (library for radiative transfer). The new parameterization considerably improves the accuracy of the retrieval of irradiances in UV bands.

1 Introduction

Radiative transfer models (RTMs) are often used to provide estimates of the UV irradiance. One of the difficulties in the computation lies in taking into account the gaseous absorption cross sections that are highly wavelength dependent (Molina and Molina, 1986). For instance, the ozone cross section changes by more than 2 orders of magnitude over the UV band [280, 400] nm. The best estimate of the UV irradiance is made by a spectrally resolved calculation of the radiative transfer for each wavelength followed by integration over the UV band. However, such spectrally detailed calculations are computationally expensive. Therefore, several methods have been proposed to reduce the number of calculations. Among them are the *k*-distribution method and the correlated-*k* approximation proposed by Kato et al. (1999). It is originally designed for providing a good estimate of the total surface solar irradiance by using 32 specific spectral intervals across the solar spectrum from 240 to 4606 nm. Hereafter, these spectral intervals are abbreviated as KBs (Kato bands). The Kato et al. method is implemented in several RTMs and is a very efficient way to speed up computations of the total surface solar irradiance. Its performance over the UV band is not very accurate when compared to detailed spectral calculations made with libRadtran (library for radiative transfer; Mayer et al., 2005) or SMARTS (Simple Model of the Atmospheric Radiative Transfer of Sunshine; Gueymard, 1995).

For a spectral interval $\Delta\lambda$ where λ is the wavelength, let $I_{0\Delta\lambda}$ and $I_{\Delta\lambda}$ denote respectively the irradiance on a horizontal plane at the top of atmosphere and at the surface; the spectral clearness index $\mathrm{KT}_{\Delta\lambda}$, also known as spectral global transmissivity of the atmosphere, spectral atmospheric transmittance, or spectral atmospheric transmission, is defined as

$$\mathrm{KT}_{\Delta\lambda} = \frac{I_{\Delta\lambda}}{I_{0\Delta\lambda}}. \quad (1)$$

Wandji Nyamsi et al. (2014) compared $KT_{\Delta\lambda}$ obtained by the correlated-k approach against that obtained by spectrally resolved computations using libRadtran and SMARTS, both for clear-sky and cloudy conditions for a set of realistic atmospheric and cloud coverage states and for each KB. They found that the Kato et al. method underestimates transmissivity in KBs #3 [283, 307] nm and #4 [307, 328] nm – covering the UV range by respectively −93 and −16 % in relative value – and exhibits relative root mean square errors (RMSEs) of 123 and 17 % in clear-sky conditions. Similar relative errors are observed for cloudy conditions.

The underestimation for these two bands can be explained by the fact that Kato et al. (1999) assume that the ozone cross section at the center wavelength in each interval represents the absorption over the whole interval. The ozone cross sections were taken from WMO (1985). Actually, the ozone cross section is strongly dependent on the wavelength in the UV region (Molina and Molina, 1986). Both KBs #3 and #4 in the UV range are large for considering only a single value of the ozone cross section.

In order to improve the potential of the Kato et al. method for estimating narrowband UV irradiances, in particular for the KBs #3 and #4, a new parameterization is proposed for the transmissivity due to the sole ozone absorption. Then, for each spectral interval, an assessment of the performance of the new parameterization in representing this transmissivity is made for a wide range of realistic cases against detailed spectral calculations. A short section describes how to implement this parameterization in the practical case of the RTM libRadtran 1.7. Finally, in each KB, the performance of the new parameterization is assessed when the direct normal, upward, downward, and global irradiances at different altitudes are computed.

2 Transmissivity due to ozone absorption

The average transmissivity $T_{o3\Delta\lambda}$ due to the sole ozone absorption for $\Delta\lambda$ can be defined by Eq. (2).

$$T_{o3\Delta\lambda} = \frac{\int_{\Delta\lambda} I_{0\lambda} e^{-k_\lambda \frac{u}{\mu_0}} d\lambda}{\int_{\Delta\lambda} I_{0\lambda} d\lambda}, \quad (2)$$

where $I_{o\lambda}$ is the spectral irradiance at the top of the atmosphere on a horizontal plane, k_λ the ozone cross section at λ, u the amount of ozone in the atmospheric column and μ_0 the cosine of the solar zenith angle.

A technique widely used for computing $T_{o3\Delta\lambda}$ is based on a discrete sum of selected exponential functions (Wiscombe and Evans, 1977):

$$T^n_{o3\Delta\lambda} = \sum_{i=1}^{n} a_i e^{-k_i\, u/\mu_0}, \quad (3)$$

where k_i are the effective ozone cross sections and a_i are the weighting coefficients obeying $\sum_{i=1}^{n} a_i = 1$.

In the Kato et al. method, only one exponential function ($n = 1$) is used for each KB to estimate the average transmissivity $T_{o3_{KB}}$:

$$T_{o3_{KB}} = e^{-k_{KB}\frac{u}{\mu_0}}. \quad (4)$$

Kato et al. (1999) have chosen the ozone cross section at the central wavelength for KB #3 and KB #4 for a temperature of 203 K: $k_{KB3} = 5.84965 \times 10^{-19}$ cm^2 and $k_{KB4} = 4.32825 \times 10^{-20}$ cm^2.

3 Effective ozone cross section

Is there a single effective ozone cross section that may represent the absorption over the whole interval? If so, this effective cross section k_{eff} is determined for each KB from the combination of Eqs. (2) and (3) with $n = 1$:

$$T_{o3_{eff}} = e^{-k_{eff}\frac{u}{\mu_0}} = \frac{1}{I_{0\Delta\lambda}} \int_{\Delta\lambda} I_{0\lambda} e^{-k_\lambda \frac{u}{\mu_0}} d\lambda. \quad (5)$$

This equation may be rewritten

$$k_{eff}\frac{u}{\mu_0} = -\ln \frac{1}{I_{0\Delta\lambda}} \int_{\Delta\lambda} I_{0\lambda} e^{-k_\lambda \frac{u}{\mu_0}} d\lambda. \quad (6)$$

Several simulations are made to study this hypothesis. The ozone cross sections are those from Molina and Molina (1986) at 226, 263 and 298 K, and the top-of-atmosphere solar spectrum of Gueymard (2004) is used. The ozone cross sections at 203 K are obtained by linear extrapolation for each wavelength (Fig. 1). Samples of 10 000 pairs (μ_0, u) were generated by a Monte Carlo technique. The random selection of the solar zenith angles follows a uniform distribution in [0°, 80°]. Similarly to what was done by Lefèvre et al. (2013) and Oumbe et al. (2014), u is computed in Dobson units as

$$u = 300\beta + 100, \quad (7)$$

where β follows the beta distribution with A parameter = 2, and B parameter = 2.

The 10 000 simulations yield a set $\mathbf{X}$ of $(\frac{u}{\mu_0})$ and a set $\mathbf{Y}$ of values $-\ln \frac{1}{I_{0\Delta\lambda}} \int_{\Delta\lambda} I_{0\lambda} e^{-k_\lambda \frac{u}{\mu_0}} d\lambda$; Eq. (6) is then

$$k_{eff}\mathbf{X} = \mathbf{Y}, \quad (8)$$

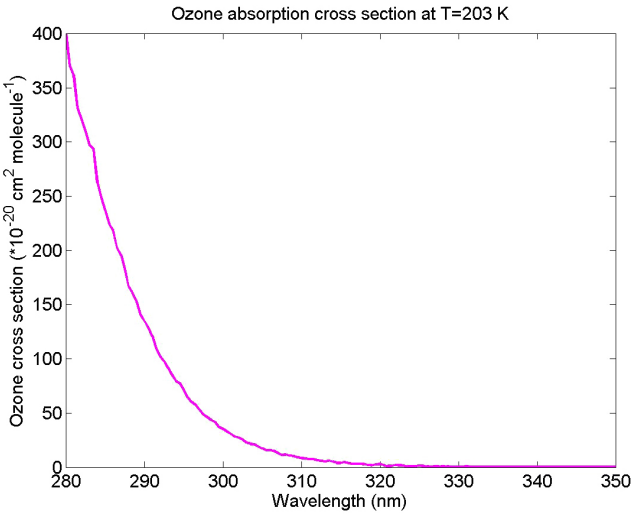

Figure 1. Ozone cross sections at 203 K as a function of the wavelength.

and k_{eff} can be found by least-square fitting technique. For KBs #3 and #4, the values obtained are respectively $k_{\text{eff3}} = 2.29 \times 10^{-19}$ cm^2 and $k_{\text{eff4}} = 2.65 \times 10^{-20}$ cm^2. The average transmissivity $T_{\text{o3}_{\text{eff}}}$ with the effective ozone cross section is then computed with Eq. (5).

Estimated transmissivities $T_{\text{o3}_{\text{KB}}}$ and $T_{\text{o3}_{\text{eff}}}$ computed with Eq. (4) and Eq. (5) using a second set of 10 000 pairs (μ_0, u) randomly selected are compared to the reference transmissivity $T_{\text{o3}\Delta\lambda}$ computed with Eq. (2) for each KB (Fig. 2). In KB #3, $T_{\text{o3}_{\text{KB}}}$ (red line) strongly underestimates $T_{\text{o3}\Delta\lambda}$, meaning that the single ozone cross section adopted by Kato et al. is too large. On the contrary, $T_{\text{o3}_{\text{eff}}}$ (blue line) exhibits a large overestimation, meaning that the efficient ozone cross section k_{eff} is too low. That may be explained by the fact that the solar radiation at the short wavelengths is completely absorbed and therefore becomes somewhat unimportant for the effective ozone cross sections. In this interval, the ozone cross section is strongly variable as shown in Fig. 1. Since k_{eff} is the optimal value reducing as much as possible the discrepancy between $T_{\text{o3}_{\text{eff}}}$ and $T_{\text{o3}\Delta\lambda}$, it may be concluded that a single effective ozone cross section may not accurately represent the absorption over the whole KB #3.

In KB #4, $T_{\text{o3}_{\text{KB}}}$ (red line) noticeably underestimates $T_{\text{o3}\Delta\lambda}$, meaning that the single ozone cross section adopted by Kato et al. is too large. $T_{\text{o3}_{\text{eff}}}$ is closer to $T_{\text{o3}\Delta\lambda}$ though it exhibits underestimation when $T_{\text{o3}\Delta\lambda} < 0.47$ and overestimation when $T_{\text{o3}\Delta\lambda} > 0.47$. Like previously stated, it may be concluded that a single effective ozone cross section may not accurately represent the absorption over the whole KB #4.

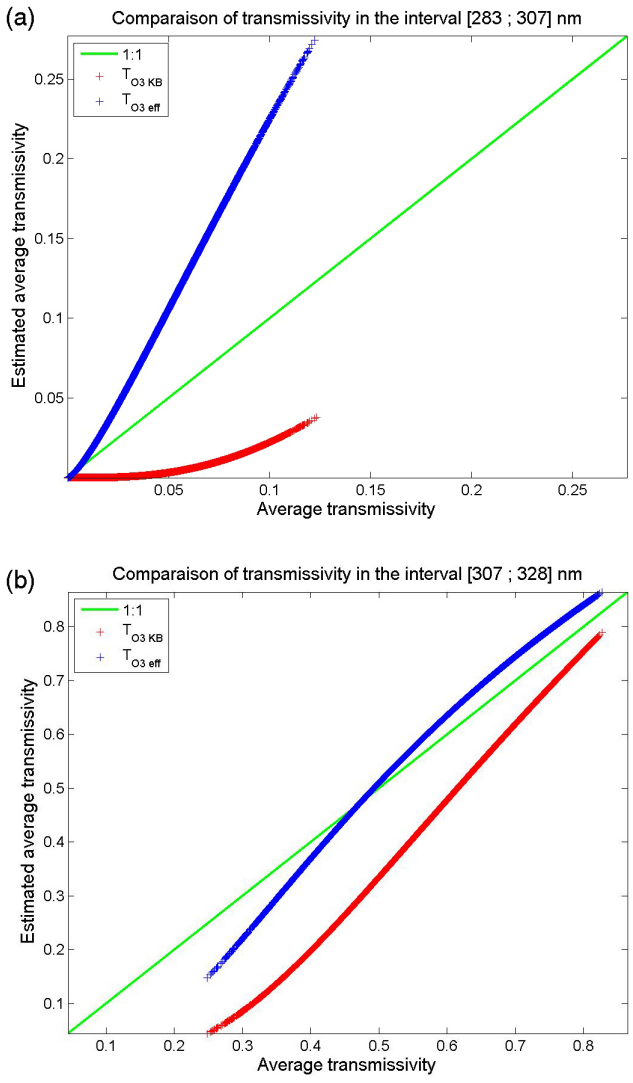

Figure 2. Scatterplot between average transmissivity $T_{\text{o3}\Delta\lambda}$ and the estimated $T_{\text{o3}_{\text{KB}}}$ (red line) and $T_{\text{o3}_{\text{eff}}}$ (blue line) for **(a)** KB #3 [283, 307] nm and **(b)** KB #4 [307, 328] nm. The identity line is in green.

4 New parameterization

The new parameterization $T_{\text{o3}_{\text{new}}}$ for computing $T_{\text{o3}\Delta\lambda}$ consists in using Eq. (3) with n greater than 1 but as small as possible to decrease the number of calculations while retaining a sufficient accuracy. n can be seen as the number of sub-intervals $\delta\lambda_i$ included in $\Delta\lambda$ for which effective ozone cross section and weighting coefficients can be defined. The greater the n, the greater the number of calculations, the more accurate the modeling of $T_{\text{o3}\Delta\lambda}$.

Many solutions are possible. No systematic scan of possible solutions in n, weight a_i and $\delta\lambda_i$ was made. This could be a further work that is computationally expensive and that requires setting up a protocol for selection of the best trade-off between accuracy and number of calculations. Here, a

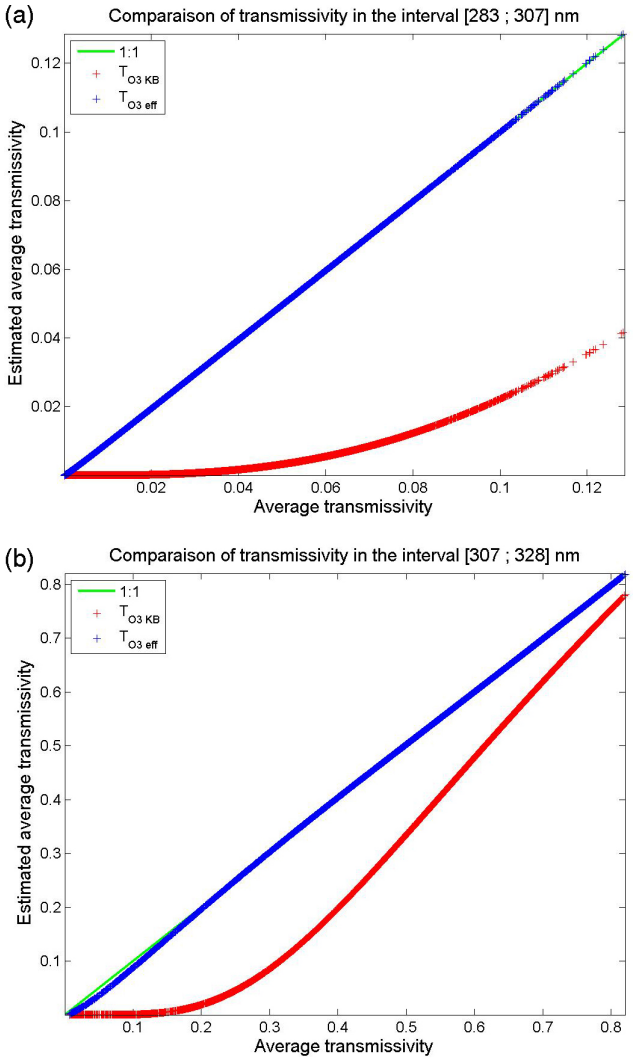

Figure 3. Scatterplot between average transmissivity $T_{o3\Delta\lambda}$ and the estimated $T_{o3_{KB}}$ (red line) and $T_{o3_{new}}$ (blue line) for **(a)** KB #3 [283, 307] nm and **(b)** KB #4 [307, 328] nm. The identity line is in green.

few tests were made with n ranging from 2 to 5. The best trade-off was found at $n = 4$. A further study was performed for $n = 4$ by adopting equal weights for the sub-intervals for both KBs #3 and #4:

$$T_{o3_{new}} = \sum_{i=1}^{4} 0.25 e^{-k_i u/\mu_0}, \quad (9)$$

where k_i is the effective ozone cross section for each of the four sub-intervals. This proposed solution is of empirical nature. Using a third set of 10 000 randomly selected pairs (μ_0, u), from which $T_{o3\Delta\lambda}$ is computed (Eq. 2), the optimal sets of four k_i and four sub-intervals $\delta\lambda_i$ minimizing the discrepancy between $T_{o3\Delta\lambda}$ and $T_{o3_{new}}$ are obtained by using the algorithm of Levenberg–Marquardt. Table 1 gives for each KB, the sub-intervals and their corresponding effective ozone cross section k_i, and weight a_i for computing $T_{o3_{new}}$. The advantage is that such parameterization is defined once for all.

Table 1. Sub-intervals, effective ozone absorption coefficient and weight in each wavelength interval for computing $T_{o3_{new}}$.

Interval $\Delta\lambda$, nm	Sub-interval $\delta\lambda_i$, nm	Effective ozone cross section k_i (10^{-19} cm^2)	Weight a_i
KB #3 283–307	283–292	11.360	0.250
	292–294	8.551	0.250
	294–301	3.877	0.250
	301–307	1.775	0.250
KB #4 307–328	307–311	0.938	0.250
	311–321	0.350	0.250
	321–323	0.153	0.250
	323–328	0.076	0.250

To assess the performance of this new parameterization, reference transmissivity $T_{o3\Delta\lambda}$ and estimated transmissivity $T_{o3_{new}}$ are computed, with respectively Eq. (2) and Eq. (9) using a fourth set of 10 000 pairs (μ_0, u) randomly selected, and compared to each other for each KB (Fig. 3). In this validation step, the random selection of the solar zenith angles follows a uniform distribution in [0°, 89°]. Statistical indicators are given in Table 2 for each KB. In general, for both KBs, the squared correlation coefficient is greater than 0.99 with very low scattering. $T_{o3_{KB}}$ (red line) is also reported in Fig. 3. The difference between $T_{o3_{KB}}$ and $T_{o3_{new}}$ is striking. In each KB, $T_{o3_{new}}$ is almost equal to $T_{o3\Delta\lambda}$ in all cases. While the mean value for $T_{o3\Delta\lambda}$ is respectively 0.0287 for KB #3 and 0.5877 for KB #4 for this data set, the maximum error in absolute value in transmissivity is respectively 0.0006 and 0.0143.

5 Practical implementation in radiative transfer model: the case of libRadtran 1.7

The file *o3.dat* in libRadtran 1.7 depicts ozone absorption. In the corresponding file, a header of seven lines describes the meanings of the following three columns. The first column contains the number of the spectral interval: KBs #1–32. The second column gives the number of quadrature points in each KB; the value is 1 in UV bands. The third column can be either the value of the single ozone cross section in each wavelength interval expressed in centimeters squared or −1 when the number of quadrature points is greater than one. In this last case, libRadtran refers to NetCDF file *cross_section.table._O3.noKB.cdf* – where *noKB* is the number of the KB – which contains the weight, the effective ozone cross section dependent of temperature and pressure.

Including the new parameterization needs two actions. Firstly, for KB #3 and KB #4, set the second column to 4 and the third column to −1. Secondly, create two NetCDF files named *cross_section.table._O3.03.cdf* and *cross_section.table._O3.04.cdf* containing for each interval

Table 2. Statistical indicators obtained by using the new parameterization for computing the transmissivity due to the sole ozone absorption in each Kato band. No. is the number of KB, R^2 is the squared correlation coefficient, mean is the mean value of the reference average transmissivity, ε is the maximum error.

No.	Mean	Bias	RMSE	rBias (%)	rRMSE (%)	R^2	ε
KB #3	0.0287	−0.0004	0.0004	−1.32	1.49	0.999	0.0006
KB #4	0.5877	−0.0005	0.0030	−0.08	0.52	0.999	0.0143

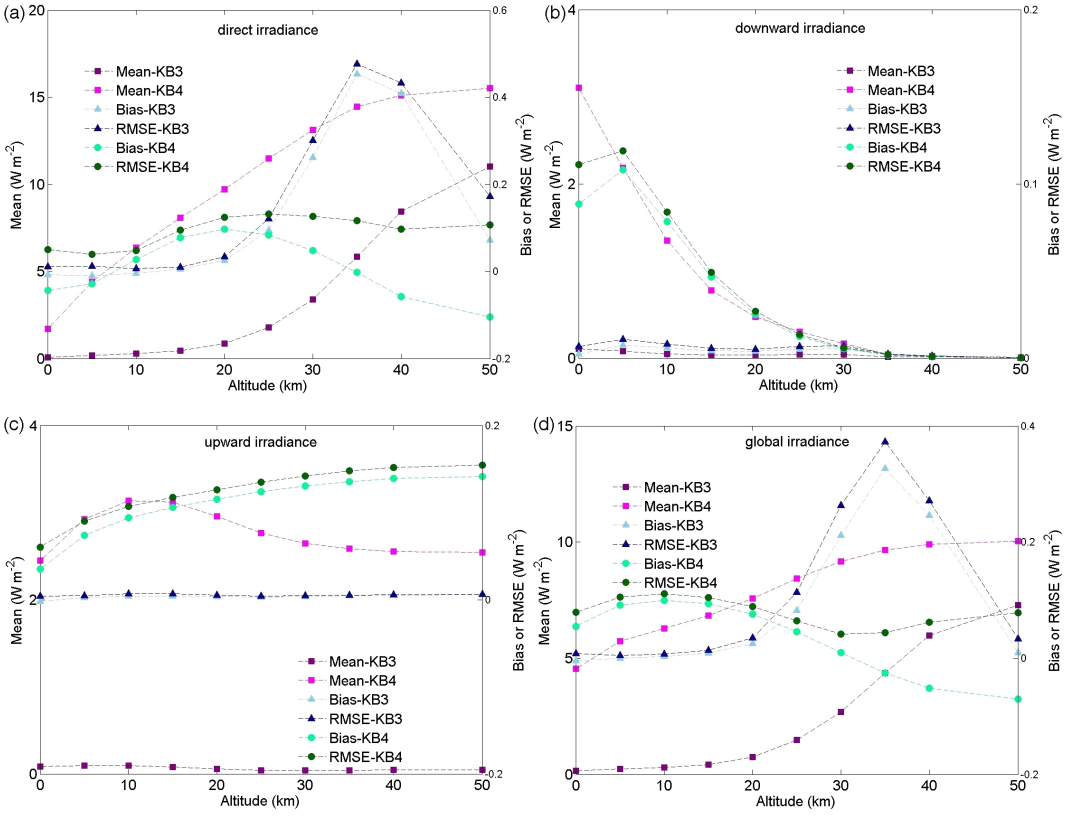

Figure 4. Mean irradiances (left vertical axis), biases and RMSEs (right vertical axis) at different altitudes in KB #3 and KB #4 for **(a)** direct normal, **(b)** downward, **(c)** upward and **(d)** global irradiance.

their corresponding weight and effective cross sections given in Table 1.

6 Performance of the new parameterization in calculating irradiances in KBs #3 and #4 in clear-sky conditions

This section presents the errors made by using the new parameterization in calculating irradiances in KBs #3 and #4. To that extent, a set of 10 000 atmospheric states have been randomly built following the marginal distribution variables described in Table 2 of Wandji Nyamsi et al. (2014), except for the solar zenith angle varying uniformly between 0 and 89°. Each atmospheric state is input to libRadtran which is run twice for KBs #3 and #4: one with detailed spectral calculations and the second with the new parameterization. The RTM libRadtran provides irradiance components that are called "direct normal", which is the irradiance received from the direction of the sun in a plane normal to the sun rays; "downward", which is the diffuse irradiance; "upward", which is the upwelling irradiance; and "global", which is the sum of the diffuse and direct irradiances, the latter being projected on a horizontal plane. Each run of libRadtran produces a set of these components at various altitudes above ground level, from 0 to 50 km, and the deviations between the irradiances produced by each run, new parameterization minus detailed spectral calculations, are computed.

The deviations are summarized by the bias, RMSE and the correlation coefficient for each altitude and in each KB (Tables 3, 4). The biases and RMSE at each altitude are summarized in Fig. 4 for both KBs. The squared correlation co-

Table 3. Statistical indicators of the performances of the new parameterization for computing the irradiances in Kato band #3 at different altitudes above ground level. Mean is the mean irradiance obtained from the detailed spectral calculations considered as reference.

Altitude (km)	KB #3							
	Direct normal irradiance ($W\,m^{-2}$)				Downward irradiance ($W\,m^{-2}$)			
	Mean	Bias	RMSE	R^2	Mean	Bias	RMSE	R^2
0	0.059	−0.008	0.011	0.999	0.108	0.002	0.007	0.999
5	0.170	−0.009	0.013	0.999	0.077	0.007	0.011	0.999
10	0.280	−0.004	0.007	0.999	0.049	0.006	0.008	0.999
15	0.454	0.005	0.010	0.999	0.034	0.004	0.006	0.999
20	0.859	0.025	0.034	0.999	0.034	0.004	0.005	0.999
25	1.784	0.094	0.121	0.999	0.041	0.005	0.007	0.999
30	3.406	0.262	0.301	0.999	0.039	0.005	0.007	0.999
35	5.832	0.453	0.476	0.999	0.015	0.002	0.002	0.996
40	8.436	0.408	0.433	0.998	0.012	0.001	0.001	0.992
50	11.024	0.072	0.178	0.998	0.005	0.000	0.000	0.999
Altitude (km)	Upward irradiance ($W\,m^{-2}$)				Global irradiance ($W\,m^{-2}$)			
	Mean	Bias	RMSE	R^2	Mean	Bias	RMSE	R^2
0	0.086	−0.002	0.004	0.999	0.162	−0.004	0.008	0.999
5	0.097	0.002	0.005	0.999	0.228	0.000	0.005	0.999
10	0.095	0.004	0.007	0.999	0.293	0.003	0.007	0.999
15	0.079	0.004	0.007	0.999	0.423	0.009	0.014	0.999
20	0.057	0.003	0.005	0.999	0.753	0.025	0.035	0.999
25	0.042	0.003	0.004	0.999	1.484	0.083	0.113	0.999
30	0.040	0.004	0.005	0.999	2.692	0.212	0.263	0.999
35	0.043	0.005	0.005	0.999	4.354	0.327	0.373	0.999
40	0.044	0.005	0.006	0.999	5.980	0.246	0.271	0.999
50	0.049	0.006	0.006	0.999	7.287	0.010	0.034	0.999

efficient is greater than 0.999, in most cases with a minimum at 0.992. This demonstrates that the new parameterization reproduces well the changes in irradiance in all cases.

The direct normal irradiance increases with altitude and exhibits negative and positive biases in both KBs #3 and #4. The bias varies as a function of the altitude. In KB #3 it reaches a minimum of $-0.009\,W\,m^{-2}$ (−5 % of the mean irradiance) at altitude of 5 km, increases with altitude up to a maximum of $0.453\,W\,m^{-2}$ (8 %) at 35 km and suddenly decreases. The RMSE follows a slightly different pattern, it decreases from $0.011\,W\,m^{-2}$ (18 % of the mean irradiance) at the surface down to a minimum $0.007\,W\,m^{-2}$ (3 %) at altitude of 10 km, then increases with altitude till a maximum of $0.476\,W\,m^{-2}$ (8 %) at 35 km and suddenly decreases. The bias and RMSE in KB #4 are less dependent on altitude. The bias is slightly negative at ground level, $-0.043\,W\,m^{-2}$ (−3 %), then increases with altitude till a maximum of $0.097\,W\,m^{-2}$ (1 %) at 20 km and gently decreases down to $-0.105\,W\,m^{-2}$ (−1 % of the mean irradiance). The RMSE is fairly constant and ranges between a minimum of $0.039\,W\,m^{-2}$ (1 %, 5 km) and a maximum of $0.132\,W\,m^{-2}$ (1 %, 25 km).

The downward irradiance decreases with altitude. The bias is positive in both KBs #3 and #4. It is fairly constant with altitude in KB #3, fluctuating between 0 and $0.007\,W\,m^{-2}$ (9 %). The bias in KB #4 decreases with altitude, from a maximum of $0.108\,W\,m^{-2}$ (5 %, 5 km) down to $0.000\,W\,m^{-2}$ at altitude of 50 km. In both KBs, the RMSE tends to decrease with altitude, from a maximum of $0.011\,W\,m^{-2}$ (14 %, 5 km), respectively $0.119\,W\,m^{-2}$ (6 %, 5 km), down to $0\,W\,m^{-2}$ at altitude of 50 km.

The upward irradiance is fairly constant with altitude in both KBs #3 and #4. The bias and the RMSE are fairly constant with altitude in KB #3, fluctuating respectively between $-0.002\,W\,m^{-2}$ (−2 %, 0 km) and $0.006\,W\,m^{-2}$ (12 %, 50 km), and between $0.004\,W\,m^{-2}$ (5 %, 0 km) and $0.007\,W\,m^{-2}$ (9 %, 15 km). The bias and RMSE in KB #4 increase with altitude. The minimum and maximum are respectively $0.035\,W\,m^{-2}$ (1 %, 0 km) and $0.141\,W\,m^{-2}$ (6 %, 50 km), and $0.006\,W\,m^{-2}$ (3 %, 0 km) and $0.155\,W\,m^{-2}$ (6 %, 50 km).

The global irradiance increases with altitude and exhibits negative and positive biases in both KBs #3 and #4. The bias varies as a function of the altitude. In KB #3, similarly to the case of the direct normal irradiance, the bias exhibits a mini-

Table 4. Statistical indicators of the performances of the new parameterization for computing the irradiances in Kato band #4 at different altitudes above ground level. Mean is the mean irradiance obtained from the detailed spectral calculations considered as reference.

Altitude (km)	KB #4							
	Direct normal irradiance ($W\,m^{-2}$)				Downward irradiance ($W\,m^{-2}$)			
	Mean	Bias	RMSE	R^2	Mean	Bias	RMSE	R^2
0	1.694	−0.043	0.050	0.999	3.105	0.088	0.111	0.999
5	4.395	−0.029	0.039	0.999	2.180	0.108	0.119	0.999
10	6.373	0.028	0.048	0.999	1.346	0.078	0.084	0.999
15	8.066	0.077	0.095	0.999	0.775	0.047	0.049	0.999
20	9.711	0.097	0.125	0.999	0.473	0.025	0.027	0.999
25	11.491	0.084	0.132	0.999	0.301	0.012	0.014	0.999
30	13.119	0.049	0.127	0.999	0.166	0.005	0.006	0.999
35	14.451	−0.002	0.117	0.999	0.042	0.002	0.002	0.999
40	15.121	−0.058	0.097	0.999	0.022	0.001	0.001	0.999
50	15.527	−0.105	0.106	0.999	0.007	0.000	0.000	0.999
Altitude (km)	Upward irradiance ($W\,m^{-2}$)				Global irradiance ($W\,m^{-2}$)			
	Mean	Bias	RMSE	R^2	Mean	Bias	RMSE	R^2
0	2.448	0.035	0.060	0.999	4.547	0.055	0.079	0.999
5	2.921	0.074	0.090	0.999	5.722	0.091	0.105	0.999
10	3.136	0.094	0.107	0.999	6.290	0.100	0.111	0.999
15	3.121	0.106	0.118	0.999	6.838	0.094	0.105	0.999
20	2.955	0.115	0.126	0.999	7.565	0.076	0.089	0.999
25	2.763	0.124	0.135	0.999	8.434	0.045	0.064	0.999
30	2.644	0.130	0.142	0.999	9.163	0.010	0.042	0.999
35	2.585	0.135	0.148	0.999	9.653	−0.025	0.044	0.999
40	2.554	0.139	0.152	0.999	9.906	−0.052	0.062	0.999
50	2.543	0.141	0.155	0.999	10.037	−0.070	0.078	0.999

mum of $-0.004\,W\,m^{-2}$ (−3 %) at the surface, then increases with altitude up to $0.327\,W\,m^{-2}$ (8 %) at 35 km and suddenly decreases down to $0.010\,W\,m^{-2}$ (0 %) at 50 km. The RMSE follows a similar trend, with a minimum of $0.005\,W\,m^{-2}$ (2 %) at altitude of 5 km, then increases up to $0.373\,W\,m^{-2}$ (9 %) at 35 km and suddenly decreases down to $0.034\,W\,m^{-2}$ (1 %) at 50 km. The situation is different in KB #4 where the bias and RMSE are less dependent with altitude. The bias is small and fluctuates between a minimum of $-0.070\,W\,m^{-2}$ (−1 %) at 50 km and a maximum of $0.100\,W\,m^{-2}$ (2 %, 10 km). The RMSE is fairly constant and ranges between a minimum of $0.042\,W\,m^{-2}$ (1 %, 30 km) and a maximum of $0.111\,W\,m^{-2}$ (2 %, 10 km).

A similar comparison was made by Wandji Nyamsi et al. (2014) with the original approach of Kato et al. (1999) but for altitudes varying between 0 and 3 km. They reported relative bias, relative RMSE and R^2 for the spectral clearness index $KT_{\Delta\lambda}$ of respectively −92 %, 123 % and 0.718 for KB #3 and −16 %, 17 % and 0.991 for KB #4. For the new parameterization, with altitudes in the range [0, 3] km, the same quantities are respectively −2 %, 4 % and 0.999 for KB #3 and −2 %, 3 % and 0.999 for KB #4. The new parameterization improves considerably the irradiances estimated in KB #3 and KB #4.

7 Conclusions

The present paper has shown the inadequacy of parameterization of the transmissivity due to the sole ozone absorption based on a single ozone cross section for the bands KB #3 [283, 307] nm and KB #4 [307, 328] nm in the *k*-distribution method and correlated-*k* approximation of Kato et al. (1999). A novel parameterization using more quadrature points better represents the transmissivity with maximum errors of respectively 0.0006 and 0.0143 for interval KBs #3 and #4. The estimates of the various components of the irradiance – direct normal, downward, upward, and global – in these Kato bands by using the new parameterization are considerably improved when compared to detailed spectral calculations. The squared correlation is greater than 0.992 in any case, and greater than 0.999 in most cases. The bias and RMSE vary with the altitude but are never greater than $0.5\,W\,m^{-2}$ for the direct normal or global in KB #3, and $0.1\,W\,m^{-2}$ in KB #4. They are smaller in KB #3 for the downward and upward irradiances ($0.01\,W\,m^{-2}$) and similar in KB #4 ($0.1\,W\,m^{-2}$).

This novel parameterization opens the way for more accurate estimates of the irradiance at the surface in the UV range and possibly in narrower spectral bands such as UV-A and UV-B.

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